On new explicit Riemannian SU(2(n+1))-holonomy metrics. ¹

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Abstract

We construct in an explicit algebraic form a family of complete noncompact Ricci-flat metrics which generalize Calabi metrics in real dimension 4(n+1) and with holonomy SU(2(n+1)).

Key words: special holonomy, Calabi metrics.

1 Introduction.

This article is concerned with an exploration of the Ricci-flat Riemannian metrics with exceptional holonomies and naturally continues a number of works [1, 2, 3, 4]. In [4] we were studying Riemannian metrics with Spin(7)-holonomy on the cones over 3-Sasakian 7-manifolds and were able to find in an explicit form a continuous family of complete noncompact 8-dimensional metrics \bar{g}_{α} depending on real parameter $0 \le \alpha \le 1$. Metric \bar{g}_0 coincides with Calabi SU(4)-holonomy metric; metric \bar{g}_1 coincides with hyperkähler Calabi metric with holonomy $Sp(2) \subset SU(4)$. Every found metric \bar{g}_{α} , $0 < \alpha < 1$ is SU(4)-holonomy metric and automatically Ricci-flat. Thus, Calabi metrics are "connected" by the obtained one-dimensional family.

However Calabi metrics (firstly appeared in [5]) are both correctly defined not only for dimension 8 but for any dimension divisible by 4. And a question of a generalization of the family constructed in [4] for higher dimensions is very natural. In this paper that question is positively resolved: for any real dimension 4(n+1) we construct in an explicit form the continuous family of metrics \bar{G}_{α} "connecting" Calabi metrics.

Theorem. The following family consists of complete Ricci-flat 4(n+1)-dimensional Riemannian metrics:

$$\bar{G}_{\alpha} = \frac{r^4 (r^4 - \alpha^4)^n}{(r^4 - \alpha^4)^{n+1} - (1 - \alpha^4)^{n+1}} dr^2 + \frac{(r^4 - \alpha^4)^{n+1} - (1 - \alpha^4)^{n+1}}{r^2 (r^4 - \alpha^4)^n} \eta_1^2 + r^2 (\eta_2^2 + \eta_3^2) + (r^2 + \alpha^2) \sum_{\beta=1}^n (\eta_{4\beta}^2 + \eta_{5\beta}^2) + (r^2 - \alpha^2) \sum_{\beta=1}^n (\eta_{6\beta}^2 + \eta_{7\beta}^2),$$

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where $0 \le \alpha \le 1$, $r \ge 1$. Metrics \bar{G}_0 and \bar{G}_1 have holonomies SU(2(n+1)) and Sp(n+1) accordingly and coincide with high-dimensional Calabi metrics found in [5]. Metrics \bar{G}_{α} for $0 < \alpha < 1$ have holonomy SU(2(n+1)) and for n = 1 coincide with the family constructed in [4]. For $0 < \alpha < 1$ metrics \bar{G}_{α} are defined on the (n+1)th tensor power of the complex line bundle over the space of complex flags in \mathbb{C}^{2n+1} and metric \bar{G}_1 is defined on $T^*\mathbb{C}P^{n+1}$.

In the next section we will explain in detail construction of the metrics \bar{G}_{α} and will prove the above theorem.

2 The Proof.

In paper [4] we explored the existence of the 8-dimensional metrics with holonomy Spin(7) of the following form

$$dt^2 + A_1(t)^2 \eta_1^2 + A_2(t)^2 \eta_2^2 + A_3(t)^2 \eta_3^2 + B(t)^2 (\eta_4^2 + \eta_5^2) + C(t)^2 (\eta_6^2 + \eta_7^2)$$

on the cone over 7-dimensional 3-Sasakian manifold M, whose 4-dimensional quaternionic Kähler orbifold \mathcal{O} possesses a Kähler structure. The form dt corresponds to the generator of the cone, the forms η_i , i=1,2,3 are the characteristic forms of the 3-Sasakian manifold and the forms η_i , $i=4,\ldots,7$ are 1-forms on the orbifold. The condition for holonomy to be contained in Spin(7) is equivalent to some system of ODE's on the functions A_i, B, C . This system was explored carefully and we were able to find the following family of solutions:

$$\bar{g}_{\alpha} = \frac{r^4(r^2 - \alpha^2)(r^2 + \alpha^2)}{r^8 - 2\alpha^4(r^4 - 1) - 1} dr^2 + \frac{r^8 - 2\alpha^4(r^4 - 1) - 1}{r^2(r^2 - \alpha^2)(r^2 + \alpha^2)} \eta_1^2 + r^2(\eta_2^2 + \eta_3^2)$$

$$+ (r^2 + \alpha^2)(\eta_4^2 + \eta_5^2) + (r^2 - \alpha^2)(\eta_6^2 + \eta_7^2),$$
(*)

where $0 \le \alpha \le 1$ and $r \ge 1$. Metrics (*) are defined on a smooth manifold if and only if M is the Aloff-Wallach space $N_{1,1} = SU(3)/S^1$.

Calabi constructed his metrics \bar{G}_0 and \bar{G}_1 on the complex bundles over the Kähler-Einstein manifolds [5]. In particular, metrics with holonomy SU(n) were constructed on the line bundles over the compact Kähler-Einstein manifolds and the hyperkähler metrics were constructed on the $T^*\mathbb{C}P^n$. Nevertheless in [5] this metrics were not written out in the explicit form.

An expression for the metric \bar{G}_0 was found in [6]:

$$\left[1 - \left(\frac{1}{\rho}\right)^{2m+2}\right]^{-1}d\rho^2 + \left[1 - \left(\frac{1}{\rho}\right)^{2m+2}\right]\rho^2(d\tau - 2A)^2 + \rho^2ds^2,\tag{1}$$

where ds^2 is a metric on a 3-dimensional Hodge and Kähler-Einstein manifold F, dA — Kähler form on F. For m=3 and $F=SU(3)/T^2$ metrics (1) and \bar{g}_0 coincide. Metric (1) is defined on the (m+1)th power of the canonical line bundle over F.

In [7] it was attempted to explore the metrics of cohomogeneity one on the $T^*\mathbb{C}P^{(n+1)}$. It's not difficult to understand the spherical subbundle in $T^*\mathbb{C}P^{(n+1)}$ fibres over the space $SU(n+2)/(U(n)\times U(1))$. On the Lie algebra su(n+1) one can choose the cobasis of the left-invariant 1-forms L_A^B such that its exterior algebra satisfies $dL_A^B = iL_A^C \wedge L_C^B$. Index A takes values in $(1,2,\beta)$ and index β takes values from 1 to n, further β is never fixed and no ambiguity appears. Obviously, $u(n) \oplus u(1)$ is a Lie subalgebra in su(n+2) and is not an exterior subalgebra. Forms $L_1^\beta = \sigma_\beta$, $L_2^\beta = \Sigma_\beta$ and $L_1^2 = \nu$ constitute a basis on the quotient $su(n+2)/(u(n) \oplus u(1))$. Then one can define the real forms: $\sigma_{1\beta} + i\sigma_{2\beta} = \sigma_\beta$, etc. Form $\lambda = L_1^1 - L_2^2$ is real by definition. In [7] the metrics of the following form

$$dt^{2} + a(t)^{2}|\sigma_{\beta}|^{2} + b(t)^{2}|\Sigma_{\beta}|^{2} + c(t)^{2}|\nu|^{2} + f(t)^{2}\lambda^{2},$$
 (2)

were considered, where the summation over the β from 1 to n is omitted. The expression for hyperkähler metric \bar{G}_1 on the $T^*\mathbb{C}P^{n+1}$ was found:

$$\frac{dr^2}{1-r^{-4}} + \frac{1-r^{-4}}{4}r^2\lambda^2 + r^2(\nu_1^2 + \nu_2^2) + \frac{r^2+1}{2}(\Sigma_{1\beta}^2 + \Sigma_{2\beta}^2) + \frac{r^2-1}{2}(\sigma_{1\beta}^2 + \sigma_{2\beta}^2). \tag{3}$$

In the case n=1 to make the notations of papers [4] and [7] agree one should put

$$\lambda = 2\eta_1, \ \nu_1 = \eta_3, \ \nu_2 = \eta_2, \ \Sigma_1 = \sqrt{2}\eta_4, \ \Sigma_2 = \sqrt{2}\eta_5, \ \sigma_1 = \sqrt{2}\eta_6, \ \sigma_2 = \sqrt{2}\eta_7.$$

In [7] also the Ricci tensor were written out. Obviously this tensor has five components: $Ric = R_0 dt^2 + R_a |\sigma_\beta|^2 + R_b |\Sigma_\beta|^2 + R_c |\nu|^2 + R_f \lambda^2$ and depends on four functions. We don't write it out here.

Notice that for any dimension n coefficients of metric (3) have the same form, and coefficients of (1) depend on n explicitly. Therefore, the requested family of metrics should depend on n explicitly and for $\alpha = 1$ should coincide with (3). We will look for metrics of the following form:

$$\frac{dr^2}{W^2} + \frac{W^2r^2}{4}\lambda^2 + r^2(\nu_1^2 + \nu_2^2) + \frac{(r^2 - \alpha^2)}{2}(\sigma_{1\beta}^2 + \sigma_{2\beta}^2) + \frac{(r^2 + \alpha^2)}{2}(\Sigma_{1\beta}^2 + \Sigma_{2\beta}^2),$$

where $W = W(r, n, \alpha)$ is an unknown function. If one will put appropriate functions into the expressions for the Ricci tensor then the components

 (R_a, R_b, R_c) will be

$$R_a = -\frac{2Q}{(r^2 - \alpha^2)^2 (r^2 + \alpha^2)}$$

$$R_b = \frac{2Q}{(r^2 + \alpha^2)^2 (r^2 - \alpha^2)}$$

$$R_c = -\frac{2Q}{r^2 (r^4 - \alpha^4)},$$

where $Q = \frac{dW}{dr}(r^5 - r\alpha^4) + 4W^2\alpha^4 + 4(n+1)(r^4 - \alpha^4 - r^4W^2)$ and $\frac{dr}{dt} = W$. This equation can be integrated without difficulties:

$$W^{2} = \frac{(r^{4} - \alpha^{4})^{n+1} + C}{r^{4}(r^{4} - \alpha^{4})^{n}},$$

where C is an integration constant. By a shift along r this constant can be fixed. One should put $C = -(1 - \alpha^4)^{n+1}$ then $r \ge 1$ and W(1) = 0. The components R_0 and R_f are the second order ODE's and automatically vanish.

Consider the following 2-form

$$\Omega = rdr \wedge \lambda + 2r^2\nu_1 \wedge \nu_2 - (r^2 + \alpha^2)\Sigma_{1\beta} \wedge \Sigma_{2\beta} + (r^2 - \alpha^2)\sigma_{1\beta} \wedge \sigma_{2\beta}.$$

Using the exterior algebra's relations from [7] one can easily verify that this form is closed and up to multiplying by $\frac{1}{2}$ is the Kähler form of metric (2). The vanishing of the Ricci tensor and the closeness of the form Ω give us the statement about holonomy.

Next we will explore the topology of the spaces where the founded metrics are defined. Here we generalize construction for higher dimensions used in [4]. Consider the complex space \mathbb{C}^{n+2} and the diagonal action of a circle S^1 on it. This action defines an equivalence class. We will designate such a class by square brackets. For example, [u], [V] where u, V are vectors or subspaces.

Consider the space $\tilde{H}=\{(u_1,u_2,V)|\ |u_1|=1,u_1\perp_{\mathbb{C}}u_2\perp_{\mathbb{C}}V\}\subset S^{2n+3}\times\mathbb{C}^{n+2}\times G_n(\mathbb{C}^{n+2}).$ Consider also the projection $\tilde{\pi}:(u_1,u_2,V)\to (u_1,V)$ from \tilde{H} to the space $\tilde{F}=\{(u_1,V)|\ |u_1|=1,u_1\perp_{\mathbb{C}}V\}.$ The spherical subbundle $\tilde{H}^1=\{(u_1,u_2,V)|\ |u_i|=1,u_1\perp_{\mathbb{C}}u_2\perp_{\mathbb{C}}V\}$ can be identified with SU(n+2)/SU(n). By using the diagonal action of the S^1 and $\tilde{\pi}$ one gets the complex line bundle $\pi:H=\tilde{H}/S^1\to F=\tilde{F}/S^1.$ The spherical subbundle in π coincides with the map $\pi^1:H^1=SU(n+2)/S[U(n)\times U(1)]\to F=SU(n+2)/T$, where

$$T = \left\{ \begin{pmatrix} z & 0 & 0 \\ 0 & \bar{z} \det \bar{A} & 0 \\ 0 & 0 & A \end{pmatrix} | z \in U(1), A \in U(n) \right\}.$$

Notice that H^1 is a 3-Sasakian manifold and coincides with Aloff-Wallach space for n=1 and π^1 is its fibration over the respective twistor space $\mathcal{Z}=F$ —the space of complex flags $\{([u],V)|\ u\in S^{2n+3}, V\in G_n(\mathbb{C}^{n+2}), u\perp_{\mathbb{C}}V\}$ in \mathbb{C}^{n+2} . It is not difficult to verify that the length with respect to the metric \bar{G}_{α} of the characteristic vector field dual to the form η_1 at the start time r=1 is equal to 2(n+1). For metric \bar{G}_{α} to be well-defined it is necessary for the circle generated by that vector field to be factorized by the discrete subgroup \mathbb{Z}_{n+1} because in the 3-Sasakian fibration π^1 there is already factorization by \mathbb{Z}_2 (look [1] for details). Thus, the metrics \bar{G}_{α} for $0 \le \alpha < 1$ are defined on the tensor power π^{n+1} . The theorem is proved.

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